# Superluminal pulse propagation in multi-level optically dressed atomic system

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Paulina Grochowska Superluminal pulse propagation

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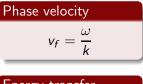
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#### Interpretation

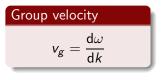


Propagation velocities Electric susceptibility Group velocity index

### Propagation velocities



Energy transfer velocity



Information transfer velocity

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Propagation velocities Electric susceptibility Group velocity index

#### It can be proved that

$$v_{g} = rac{c}{1+rac{\omega}{2}rac{\mathrm{d}\chi'(\omega)}{\mathrm{d}\omega}|_{\omega=\omega_{0}}} = rac{c}{n_{g}}$$

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Propagation velocities Electric susceptibility Group velocity index

## Analysis of group velocity index

The most interesting case is for  $n_g \in (-\infty, 1)$ .

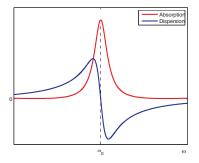
In this case we obtain a <u>negative</u> group velocity or group velocity greater than c.

To have  $n_g \in (-\infty, 1)$  we should find such an atomic system in which the dispersion is normal while absorption (or gain) being not too strong.

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Propagation velocities Electric susceptibility Group velocity index

## The Electric susceptibility

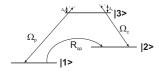


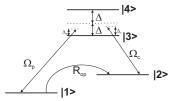
**Rysunek**: A typical plot of  $\chi(\omega)$  (real part - dispersion and imaginary part - absorption)

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Atomic systems Hamiltonian The electric susceptibility

#### Atomic systems





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Atomic systems Hamiltonian The electric susceptibility

## Hamiltonian

#### Hamiltonian of the system

$$\mathcal{H} = egin{bmatrix} \Delta_{
ho} & 0 & \Omega_{
ho3} & \Omega_{
ho4} \ 0 & \Delta_c & \Omega_{c3} & \Omega_{c4} \ \Omega_{
ho3}^* & \Omega_{c3}^* & \Delta & 0 \ \Omega_{
ho4}^* & \Omega_{c4}^* & 0 & -\Delta \end{bmatrix}$$

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Atomic systems Hamiltonian The electric susceptibility

## Hamiltonian

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ho4}^* & \Omega_{c4}^* & 0 & -\Delta \end{bmatrix}$$

Von Neumann's equation:

$$i\hbar \frac{\mathrm{d}\rho}{\mathrm{d}t} = [\mathcal{H},\rho] + \Lambda$$

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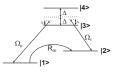
Atomic systems Hamiltonian The electric susceptibility

## The electric susceptibility in atomic systems

The electric susceptibility in atomic systems can be written as:

$$\chi = \frac{N\left(|d_{13}|^2 \rho_{13} + |d_{14}|^2 \rho_{14}\right)}{\varepsilon_0 \hbar \Omega_p}$$

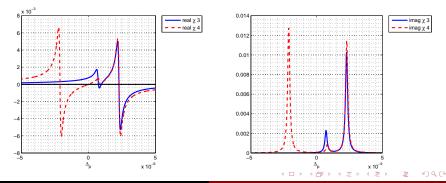
By solving the system of equations of motion for density matrix and putting the solutions to the expression above we obtain graphs for dispersion and absorption depending on following parameters:  $\Delta, \Delta_c, R_{op}, \Omega_c, \Omega_p$  as a function of  $\Delta_p$ .



without optical pump With optical pump  $v_g \in (-\infty, 0)$  $v_g \in (c, \infty)$ Second derivative approximation Third derivative approximation

#### Results - without optical pump

 $\Delta = 2 \cdot 10^{-9}, \Delta_c = 1 \cdot 10^{-9}, E_c = 6 \cdot 10^{-10}, E_p = 3 \cdot 10^{-12}$ 

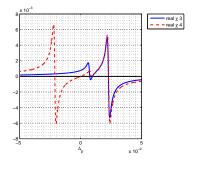


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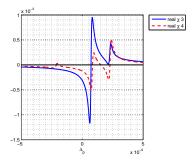
#### Superluminal pulse propagation

without optical pump With optical pump  $v_g \in (-\infty, 0)$   $v_g \in (c, \infty)$ Second derivative approximation Third derivative approximation

#### Optical pump influence on dispersion



Rysunek: Without optical pump



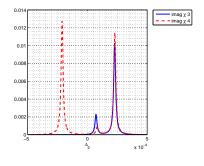
**Rysunek:** With optical pump  $R_{op} = 2 \cdot 10^{-9}$ 

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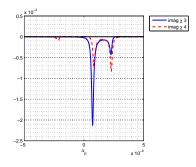
Paulina Grochowska Superluminal pulse propagation

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 $\begin{array}{c} \mbox{Motivation} & \mbox{without optical pump} \\ \mbox{Calculations} & \mbox{With optical pump} \\ \mbox{Weth optical pump } \\ \mbox{Weth opt$ 

$$v_g \in (-\infty, 0)$$

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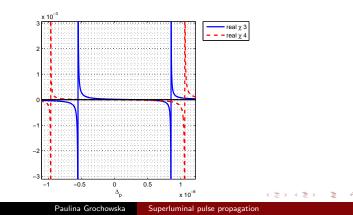
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without optical pump With optical pump  $\mathbf{vg} \in (-\infty, \mathbf{0})$  $\mathbf{vg} \in (c, \infty)$ Second derivative approximation Third derivative approximation

$$\frac{\mathsf{Results} - v_g \in (-\infty, 0)}{\mathsf{Results} - v_g \in (-\infty, 0)}$$

Analysis of dispersion

 $\Delta = 2 \cdot 10^{-9}, \Delta_c = 1 \cdot 10^{-9}, E_c = 7 \cdot 10^{-9}, E_p = 3 \cdot 10^{-12}, \textit{Rop} = 2 \cdot 10^{-11}$ 

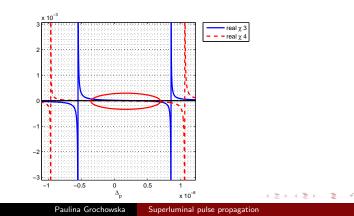


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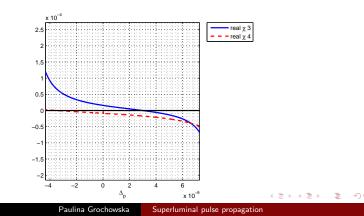


without optical pump With optical pump  $\mathbf{v}_{\mathbf{g}} \in (-\infty, \mathbf{0})$  $\mathbf{v}_{\mathbf{g}} \in (c, \infty)$ Second derivative approximation Third derivative approximation

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Analysis of dispersion

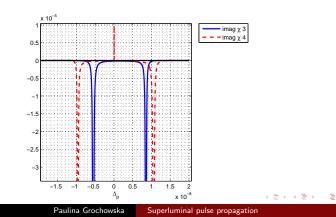
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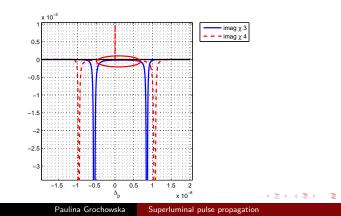
Analysis of absorption  $\Delta = 2 \cdot 10^{-9}, \Delta_c = 1 \cdot 10^{-9}, E_c = 7 \cdot 10^{-9}, E_p = 3 \cdot 10^{-12}, Rop = 2 \cdot 10^{-11}$ 



without optical pump With optical pump  $v_g \in (-\infty, 0)$  $v_g \in (c, \infty)$ Second derivative approximation Third derivative approximation

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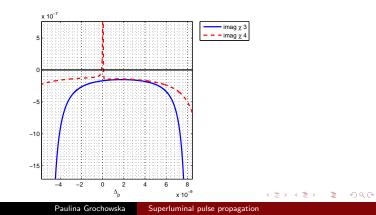


without optical pump With optical pump  $\mathbf{v}_{\mathbf{g}} \in (-\infty, \mathbf{0})$  $\mathbf{v}_{\mathbf{g}} \in (c, \infty)$ Second derivative approximation Third derivative approximation

$$\frac{\mathsf{Results} - v_g \in (-\infty, 0)}{\mathsf{Results}}$$

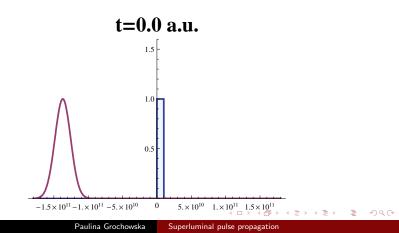
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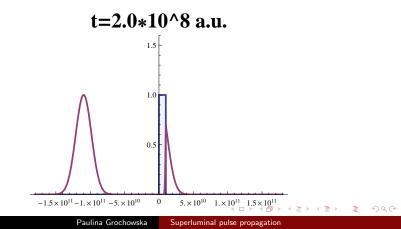
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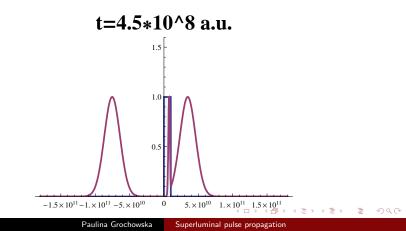
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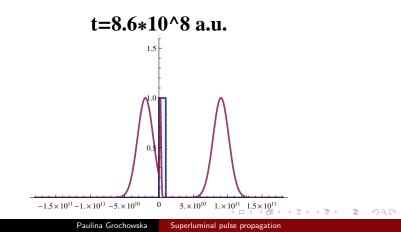
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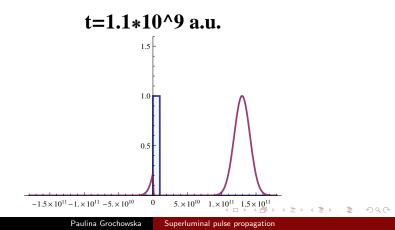
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without optical pump With optical pump  $v_g \in (-\infty, 0)$  $v_g \in (c, \infty)$ Second derivative approximation Third derivative approximation

$$v_g \in (c,\infty)$$

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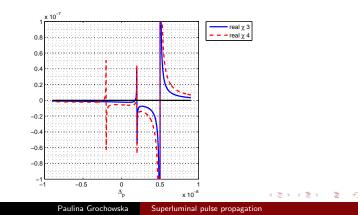
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without optical pump With optical pump  $v_g \in (-\infty, 0)$  $v_g \in (c, \infty)$ Second derivative approximation Third derivative approximation

## Results - $v_g \in (c,\infty)$

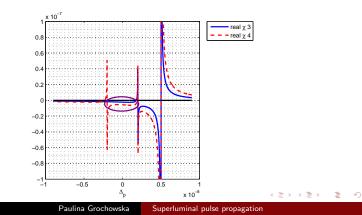
Analysis of dispersion



without optical pump With optical pump  $v_g \in (-\infty, 0)$  $v_g \in (c, \infty)$ Second derivative approximation Third derivative approximation

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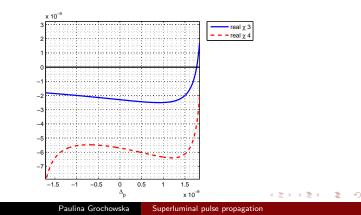
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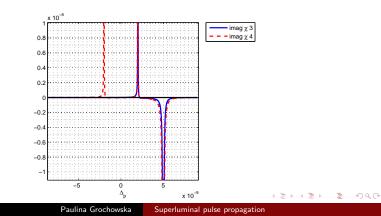
Analysis of dispersion



without optical pump With optical pump  $v_g \in (-\infty, 0)$  $v_g \in (c, \infty)$ Second derivative approximation Third derivative approximation

## Results - $v_g \in (c,\infty)$

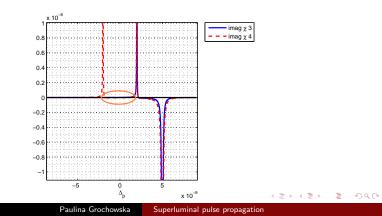
Analysis of absorption



without optical pump With optical pump  $v_{g} \in (-\infty, 0)$  $v_{g} \in (c, \infty)$ Second derivative approximation Third derivative approximation

## Results - $v_g \in (c,\infty)$

Analysis of absorption

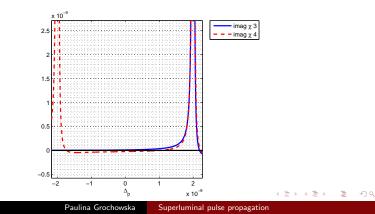


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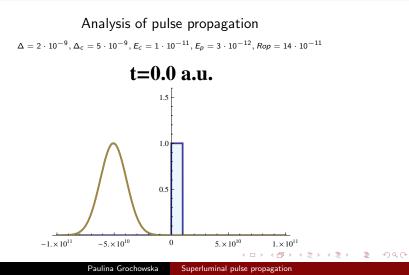
## $\frac{|\mathsf{Results} - \textit{v}_g \in (c,\infty)|}{|\mathsf{Results} - \mathsf{v}_g \in (c,\infty)|}$

Analysis of absorption

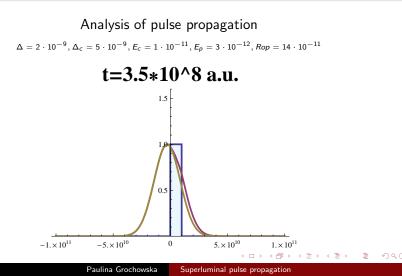
$$\Delta = 2 \cdot 10^{-9}, \Delta_c = 5 \cdot 10^{-9}, E_c = 1 \cdot 10^{-11}, E_p = 3 \cdot 10^{-12}, Rop = 14 \cdot 10^{-11}, n_g = 0.95 \cdot 10^{-11}, E_p = 10^{-11}, E_p$$



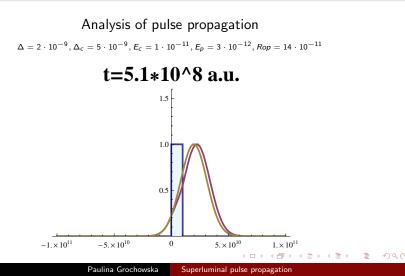
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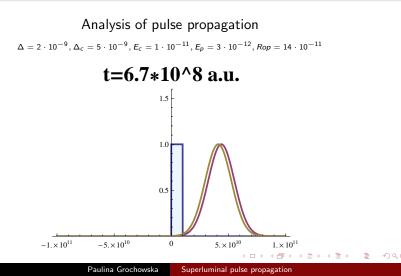
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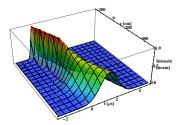


without optical pump With optical pump  $v_g \in (-\infty, 0)$  $v_g \in (c, \infty)$ Second derivative approximation Third derivative approximation

### Second derivative approximation

An example of an impulse propagating through a sample with a negative group velocity for electric susceptibility approximated by:

$$\tilde{\chi}(\Delta_{p}) = \chi(0) + \frac{d\chi(0)}{d\omega} (\Delta_{p}) + \frac{1}{2} \frac{d^{2}\chi(0)}{d\Delta_{p}^{2}} (\Delta_{p})^{2}$$



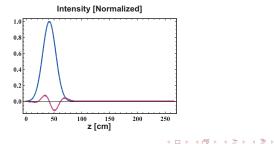
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without optical pump With optical pump  $v_g \in (-\infty, 0)$  $v_g \in (c, \infty)$ Second derivative approximation Third derivative approximation

#### Third derivative approximation

A pulse propagating in a subluminal regime. The Brillouin precursor magnified by the factor of 5000 is shown. The precursor was obtained for electric susceptibility approximated by:

$$\tilde{\chi}\left(\Delta_{\rho}\right) = \chi(0) + \frac{\mathrm{d}\chi(0)}{\mathrm{d}\omega} \left(\Delta_{\rho}\right) + \frac{1}{2} \frac{\mathrm{d}^{2}\chi(0)}{\mathrm{d}\Delta_{\rho}^{2}} \left(\Delta_{\rho}\right)^{2} + \frac{1}{6} \frac{\mathrm{d}^{3}\chi'(0)}{\mathrm{d}\Delta_{\rho}^{3}} \left(\Delta_{\rho}\right)^{3}$$



### Interpretation on superluminal propagation

Problems:

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## Interpretation on superluminal propagation

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# Interpretation on superluminal propagation

Problems:

- Conservation of energy
  - There is an exchange of energy between the medium and the front and back parts of the pulse, leading to the pulse advancement in presence of the optical pump.
     G. Diener, Phys. Lett. A 235 (1997) 118-124

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- Effects are observed only for analytical pulses.

#### Summary

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**1** 
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 $\bullet~$  In high detuning regimes  $\Delta_c > 3\cdot 10^{-9}$  for  $\Delta = 2\cdot 10^{-9}$ 

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There have been demonstrated two kinds of superluminal pulses (in terms of group velocity):

$$\textcircled{0} \ v_g > c \ \text{for} \ n_g \in (0,1)$$

- In high detuning regimes  $\Delta_c > 3 \cdot 10^{-9}$  for  $\Delta = 2 \cdot 10^{-9}$
- With relatively weak coupling field  $E_c = 1 \cdot 10^{-11}$  for  $E_{\rho} = 3 \cdot 10^{-10}$

**2** Negative  $v_g$  for  $n_g \in (-\infty, -1)$ 

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There have been demonstrated two kinds of superluminal pulses (in terms of group velocity):

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**2** Negative  $v_g$  for  $n_g \in (-\infty, -1)$ 

• In both high and weakly detuning regimes

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- **2** Negative  $v_g$  for  $n_g \in (-\infty, -1)$ 
  - In both high and weakly detuning regimes
  - For relatively strong coupling field  $E_c = 3 \cdot 10^{-9}$  for  $E_p = 3 \cdot 10^{-10}$
  - In a wide range of frequency

# Summary

There have been demonstrated two kinds of superluminal pulses (in terms of group velocity):

**0** 
$$v_g > c$$
 for  $n_g \in (0, 1)$ 

- $\bullet~$  In high detuning regimes  $\Delta_c > 3\cdot 10^{-9}$  for  $\Delta = 2\cdot 10^{-9}$
- With relatively weak coupling field  $E_c = 1 \cdot 10^{-11}$  for  $E_p = 3 \cdot 10^{-10}$
- In a narrow range of frequency
- 2 Negative  $v_g$  for  $n_g \in (-\infty, -1)$ 
  - In both high and weakly detuning regimes
  - For relatively strong coupling field  $E_c = 3 \cdot 10^{-9}$  for  $E_p = 3 \cdot 10^{-10}$
  - In a wide range of frequency

In both cases we can control the value of the group velocity index by changing the amplitude of coupling field or incoherent optical pump.



All the results have been obtained in cooperation with:

- prof. A. Raczyński<sup>1</sup>
- prof. J. Zaremba<sup>1</sup>
- dr S. Zielińska-Kaniasty<sup>2</sup>

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#### Thank you for your attention

Paulina Grochowska Superluminal pulse propagation

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