

Entanglement and sensitivity in precision measurements with a fluctuating number of particles

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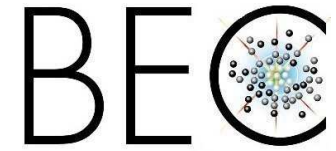
(to appear in PRL)

Torun, August 31, 2010



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Introduction



Aims of the project

Extend results on phase estimation with linear interferometers regarding

- Shot-noise limit vs. Spin-Squeezing and Entanglement
- Heisenberg limit

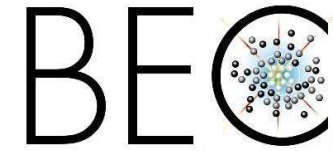
to probe states with a non-fixed number of particles N

Motivation

- Fundamental extension of the theory
- Recent experimental works on Spin-Squeezing
 - work with a fluctuating number of atoms
 - use theory developed for fixed N ***Is that justified?***

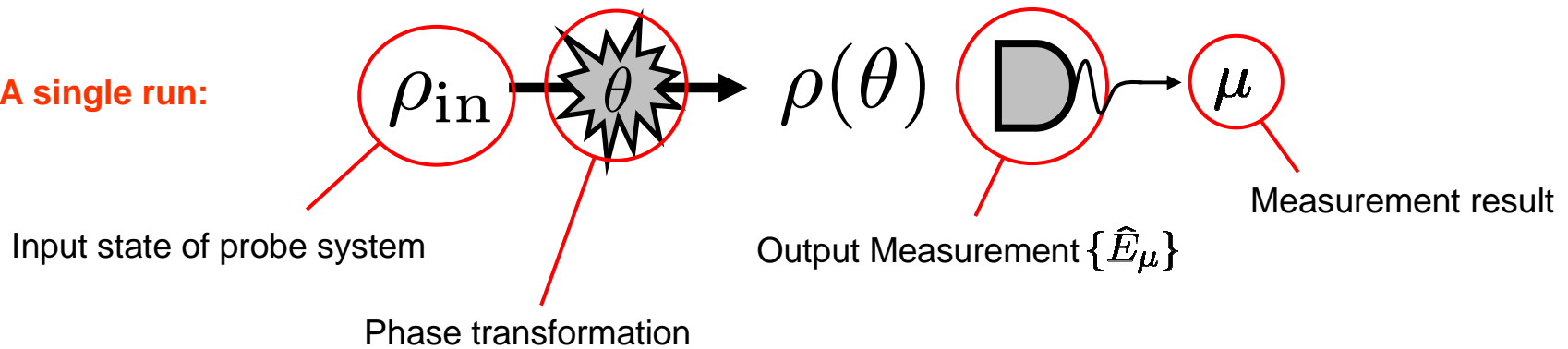
Jessen (cold atoms) [PRL 2007](#)
Oberthaler (BEC) [Nature 2008 and 2010](#)
Polzik (cold atoms) [PNAS USA 2009](#)
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Treutlein (BEC on atom chip) [Nature 2010](#)
...

Introduction



Phase Estimation

1) A single run:



2) m runs → $\{ \mu_i \}_{i=1}^m$ → Phase estimator $\hat{\theta}(\{ \mu_i \})$

3) Cramer-Rao bound

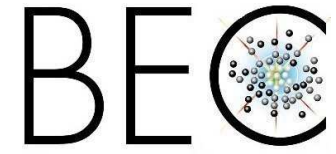
$$\Delta \hat{\theta} \geq \frac{1}{\sqrt{m F}}$$

Fisher information F

$\left[\begin{array}{l} \text{input state } \rho_{\text{in}} \\ \text{phase operation} \\ \text{output measurement} \end{array} \right]$

(For unbiased estimators $\langle \hat{\theta} \rangle = \theta$)

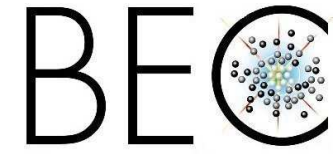
Introduction



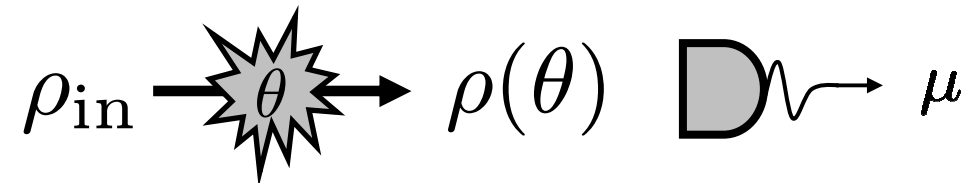
In the following slides

fixed N

Introduction



Linear Interferometers



two-mode approximation interferometer acts only in a two-level subspace \longrightarrow "spins"

unitary linear interferometer $\rho(\theta) = e^{-i\hat{H}\theta} \rho_{\text{in}} e^{i\hat{H}\theta}$, $\hat{H} = \frac{1}{2} \sum_{i=1}^N \hat{\sigma}_{\vec{n}_i}$

Shot-Noise Limit and Entanglement

For separable input states $\rho_{\text{sep}} = \sum_k p_k |\psi_{\text{sep}}^{(k)}\rangle \langle \psi_{\text{sep}}^{(k)}|$

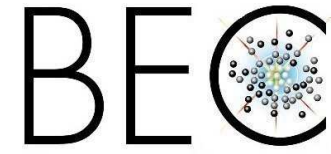
$$|\psi_{\text{sep}}\rangle = |\psi_1\rangle \otimes |\psi_2\rangle \otimes \dots \otimes |\psi_N\rangle$$

state of particle 1
with two degrees of freedom

and linear interferometers: $F \leq N$ Pezze and Smerzi, PRL 09

$\Delta\hat{\theta} \geq \frac{1}{\sqrt{mF}} \geq \frac{1}{\sqrt{mN}}$ **Entanglement is necessary for Sub Shot-Noise Interferometry**

Introduction



Spin Squeezing

Kitagawa and Ueda, PRA 1993; Wineland et al., PRA 1994

Collective Spin Operators $\hat{J}_k = \frac{1}{2} \sum_{l=1}^N \hat{\sigma}_k^{(l)}$, $k = x, y, z$, $[\hat{J}_k, \hat{J}_j] = i\epsilon_{kjr} \hat{J}_r$

Spin Squeezing Parameter $\xi = \frac{N \Delta \hat{J}_z}{|\langle \hat{J}_x \rangle|}$ if $\xi < 1$

- state useful for Sub Shot-Noise Interferometry
- state of the atoms is entangled *Soerensen et al., Nature 2001*

Fisher vs Spin-Squeezing: $\chi^2 \equiv \frac{N}{F} \leq \xi^2 \longrightarrow$ Spin-Squeezing implies $F > N$

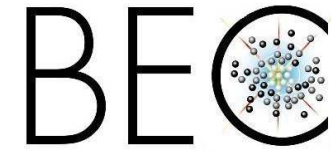
Pezze and Smerzi, PRL 09

Heisenberg limit

$$F \leq N^2 \longrightarrow \Delta \hat{\theta} \geq \frac{1}{\sqrt{mN^2}} \quad \text{Giovannetti et al., PRL 06}$$

$F = N^2$ saturated by Cat/NOON/GHZ state

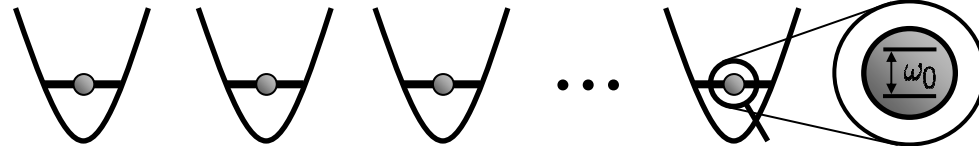
Introduction



Wineland et al., PRA 1994

Illustration: Ramsey spectroscopy in an ion trap

System: N locally addressable ions
with two internal states $|z_{\pm}\rangle$



Aim: Measure transition frequency ω_0 between internal states

- Tools:**
- Manipulate internal states with classical radiation
 - Measurement of populations in the states $|z_{\pm}\rangle$

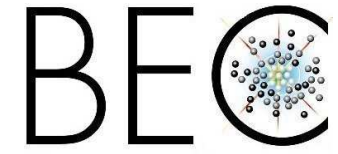
(Modified) **Ramsey scheme:**

1) Initialization: Prepare all atoms in $|z_{-}\rangle$

2) Sequence: $\rho_{\text{in}} \xrightarrow{e^{-i\hat{J}_y\omega_0 T}} \rho(\omega_0 T) \quad \hat{J}_z \mathcal{N}$

Linear Interferometer with
 $\hat{H} = \hat{J}_y$ and $\theta = \omega_0 T$
and addressable particles
 Separable input state
 → **Shot Noise Limited**

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 - work with a fluctuating number of atoms
 - use simple extension of the theory developed for fixed N

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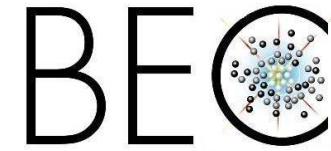
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...

Is that justified?

Non-fixed N



Separability and Entanglement



What kind of states are possible?

- only incoherent mixtures $\rho_{\text{inc}} = \sum_N Q_N \rho^{(N)}$
- or states with coherences $|\psi_{\text{coh}}\rangle = \sum_N \sqrt{Q_N} |\psi^{(N)}\rangle$?

Super Selection Rules (SSRs) generally forbid the creation and detection of such coherences ...
 ... but this may be possible if suitable *Reference Frames* can be established.

Wick et al., PR 52, Moelmer, PRA 97, Bartlett et al., RMP 07, ...

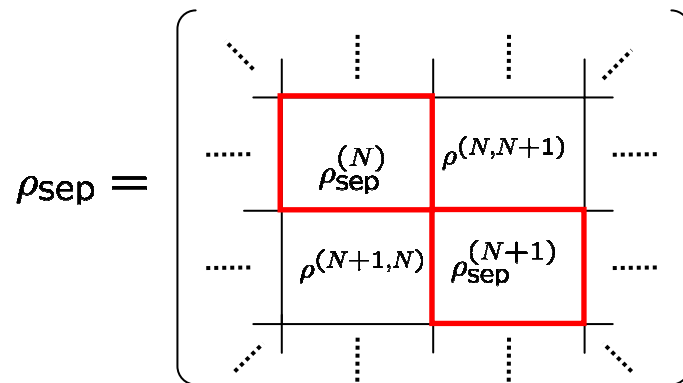
① SSR: $\rho_{\text{sep}} = \sum_N Q_N \rho_{\text{sep}}^{(N)}$

recall:

$$\rho_{\text{sep}}^{(N)} = \sum_k p_k |\psi_{\text{sep}}^{(k,N)}\rangle \langle \psi_{\text{sep}}^{(k,N)}|,$$

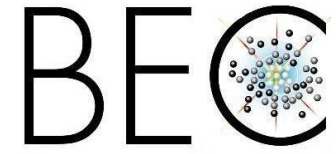
$$|\psi_{\text{sep}}^{(N)}\rangle = |\psi_1\rangle \otimes |\psi_2\rangle \otimes \dots \otimes |\psi_N\rangle$$

② COH:



Non-separable states are **entangled**

Non-fixed N

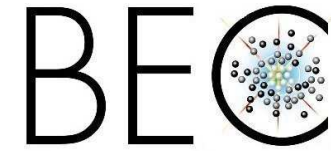


Collective Spin Operators

$$\hat{J}_i \rightarrow \begin{pmatrix} \ddots & & & \\ \dots & \hat{J}_i^{(N)} & 0 & \dots \\ \dots & 0 & \hat{J}_i^{(N+1)} & \dots \\ \dots & & & \ddots \end{pmatrix}$$

$$\hat{J}_i^{(N)} = \frac{1}{2} \sum_{j=1}^N \hat{\sigma}_i^{(j)}$$

Non-fixed N - Results



	fixed- N	non-fixed- N : SSR	non-fixed- N : COH
Shot-Noise limit = Minimal $\Delta\theta$ with separable states	$\Delta\theta \geq \frac{1}{\sqrt{mN}}$ $\chi^2 = \frac{N}{F}$	$\Delta\theta \geq \frac{1}{\sqrt{m\langle\hat{N}\rangle}}$ $\chi^2 = \frac{\langle\hat{N}\rangle}{F}$	
Spin-Squeezing	$\chi \leq \xi \equiv \frac{N\Delta\hat{J}_z}{ \langle\hat{J}_x\rangle }$	$\chi \leq \xi \equiv \frac{\langle\hat{N}\rangle\Delta\hat{J}_z}{ \langle\hat{J}_x\rangle }$	
Heisenberg limit	$\Delta\theta \geq \frac{1}{\sqrt{mN^2}}$	$\Delta\theta \geq \max\left[\frac{1}{\sqrt{m\langle\hat{N}^2\rangle}}, \frac{1}{m\langle\hat{N}\rangle}\right]$	$\Delta\theta \geq \frac{1}{\sqrt{m\langle\hat{N}^2\rangle}}$ (but $\Delta\theta \geq \frac{1}{m\langle\hat{N}\rangle}$ in central limit)

Details in: P. Hyllus, L. Pezze, and A. Smerzi, arXiv:1003.0649 (to appear in PRL)

Multiparticle entanglement, Spin Squeezing, and Phase Estimation → arXiv:1006.4366 and arXiv:1006.4368

Thank you very much for your attention!