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Macroscopic quantum states of light generated by quantum cloners

Motivation

Fundamental interest: quantum-classical border

loophole-free Bell tests

Macroscopic nature makes them <u>always</u> detectable with classical detectors

Applications:



quantum metrology (super-resolution)

quantum information technology (efficient coupling between light and matter)

Multiphoton Polarization Entanglement



Population on average $4sh^2g$ of photons

F. De Martini et al, PRL 100, 253601 (2008)

Multiphoton Polarization Entanglement

Equatorial states of polarization

$$\begin{aligned} |1_{\varphi}\rangle &= 1/\sqrt{2}(|1_H, 0_V\rangle + e^{i\varphi}|0_H, 1_V\rangle) \\ |1_{\varphi^{\perp}}\rangle &= 1/\sqrt{2}(|1_H, 0_V\rangle - e^{i\varphi}|0_H, 1_V\rangle) \end{aligned}$$



PDC singlet

$$1/\sqrt{2}(|1_{\varphi}\rangle_{A}|1_{\varphi^{\perp}}\rangle_{B} - |1_{\varphi^{\perp}}\rangle_{A}|1_{\varphi}\rangle_{B})$$

Quantum amplification

 $|1_{\varphi}\rangle \rightarrow |\Phi\rangle \quad |1_{\varphi^{\perp}}\rangle \rightarrow |\Phi_{\perp}\rangle$

Macroscopic Quantum Properties

Polarization and photon number properties

$$\begin{split} |\Phi\rangle \\ (1,3,5,\dots) \text{ photons in } \varphi \\ (0,2,4,\dots) \text{ photons in } \varphi^{\perp} \end{split}$$



$$\bar{m}_{\varphi} = 3m + 1$$
$$\bar{m}_{\varphi^{\perp}} = m$$

 $|\Phi_{\perp}\rangle$ (1,3,5,...) photons in φ^{\perp} (0,2,4,...) photons in φ

 $\bar{m}_{\varphi^{\perp}} = 3m + 1$ $\bar{m}_{\varphi} = m$

Population $\bar{m} = 10^5$ for g = 6

"Micro-macro" singlet

$$1/\sqrt{2}\left(|1_{\varphi}\rangle_{A}|\Phi_{\perp}\rangle_{B}-|1_{\varphi^{\perp}}\rangle_{A}|\Phi\rangle_{B}\right)$$

"Macro-macro" singlet

 $1/\sqrt{2}\left(|\Phi\rangle_A|\Phi_{\perp}\rangle_B - |\Phi_{\perp}\rangle_A|\Phi\rangle_B\right)$



Singlets are useful if $|\Phi
angle$ and $|\Phi_{\perp}
angle$ are distinguishable

Photon Number Distribution: detector point of view

Effective overlap ~ 10^{-1}

Photon distributions are continuous

$$p_{\Phi}(k,l) = |\langle k, l | \Phi \rangle|^2$$
$$p_{\Phi^{\perp}}(k,l) = p_{\Phi}(l,k)$$

Distinguishability v = 0.64



How to improve distinguishability of macroscopic states?

Filtering – Quantum State Engineering

POVM (positive operator valued measure) describe measurement induced quantum operations

preserve quantum macroscopic character

significantly improve distinguishability with inefficient detection

are crucial for any quantum protocol, Bell inequality violation, quantum cryptography and quantum metrology

Modulus Difference Filter

Quantum scissors

$$\mathcal{P}_{D_{th}} = \sum_{\substack{k,l=0\\|k-l| \ge D_{th}}}^{\infty} |k,l\rangle \langle k,l|$$

POVM measurement

MDF filters both macroqubits fairly



MDF has binary output (exact value of the modulus is not measured)

engineering of quantum superposition without collapsing on a particular Fock state







Filtering in Presence of Losses

 $\bar{m} = 10$

50%





Distinguishability improvement is noticible for detectors measuring $\pm 150~\rm{photons}$

Physical Implementation of Modulus Difference Filter

"Tap" measurement, feed forward technique, and Hong-Ou-Mandel effect



Hong-Ou-Mandel Effect

Interference of two non-zero Fock states on BS





Multiphoton Hong-Ou-Mandel Effect

PBS and output: (H, V) basis Input: $(H + e^{i\varphi}V, H - e^{i\varphi}V)$ basis



Detected state

$$|K,L
angle_{HV}$$
 $\sigma = K+l$

The input state is given by BS operation

$$|\Psi\rangle_{in} = \mathcal{N}\sum_{\delta=-\sigma}^{\sigma} A^{\sigma}\left(\frac{\sigma+\delta}{2}, \frac{\sigma-\delta}{2}\right) \left|\frac{\sigma+\delta}{2}, \frac{\sigma-\delta}{2}\right\rangle_{r,r^{\perp}}$$

Multiphoton Hong-Ou-Mandel Effect

 $\sigma = 200$



Bell Inequality Test for Micro-Macro Singlet with Preselection

 $\frac{1}{\sqrt{2}}\left(|1_{\varphi}\rangle_{A}|\Phi_{\perp}\rangle_{B}-|1_{\varphi^{\perp}}\rangle_{A}|\Phi\rangle_{B}\right)$



Bell Inequality Test for Micro-Macro Singlet with Preselection



+1 -1

BS $r \ll t$

 $\begin{array}{c|c} & \langle \mathbf{u}, \mathbf{u} \\ \hline & |\mathbf{1}_{\varphi} \rangle \\ \hline & |\mathbf{1}_{\varphi^{\perp}} \rangle \\ \hline & \varphi \\ \hline & & |\Phi_{\perp}\rangle \end{array} \begin{array}{c} |\Phi\rangle \\ \hline & |\Phi_{\perp}\rangle \end{array}$

 β, β'

Correlation function

 $E(\alpha,\beta) = \langle \mathcal{O}^{(m)} \otimes \mathcal{O}^{(M)} \rangle$

CHSH-Bell inequality $B = E(\alpha, \beta) + E(\alpha', \beta) + E(\alpha, \beta') - E(\alpha', \beta')$ Difference operator is not rotationally invariant

If distinguishability increases in the source basis $\varphi = \pi/2$, it decreases in the orthogonal basis $\varphi = 0$



For every preselection basis θ there are rotations $\alpha,\,\beta$ for which B>2

Analytical result

Numerical result



Maximal value of B is independent of preselection basis θ

Numerical Results g = 1.1 (7 photons on average)

 $D_{th} = 7, \ B = \pm 1.91$







Conclusions

We proposed Modulus Intensity Difference filter

POVM measurement (preserves quantum superposition) Based on Hong-Ou-Mandel effect Works for super-Poissonian input statistics

Applied to quantum macroscopic states of light

Increases distinguishability, also in presence of high losses Increases sensitivity of phase estimation Hope for a loophole-free Bell inequality test





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Thank You!

