

# Quantum simulation of an extra dimension



**Alessio Celi**

based on PRL 108, 133001 (2012), with O. Boada, J.I. Latorre, M. Lewenstein,

**Quantum Technologies Conference III**

# Outline:

- Idea of Quantum Simulator ([Feynman](#))

# Outline:

- Idea of Quantum Simulator ([Feynman](#))
- Cold Atoms

# Outline:

- Idea of Quantum Simulator ([Feynman](#))
- Cold Atoms
  - **Optical lattice**: Hubbard model with internal degrees of freedom

# Outline:

- Idea of Quantum Simulator ([Feynman](#))
  - Cold Atoms
    - **Optical lattice**: Hubbard model with internal degrees of freedom
- “Old” Ideas (and applications):
- Synthetic gauge field: [Jaksch and Zoller proposal and extensions](#)

# Outline:

- Idea of Quantum Simulator ([Feynman](#))
- Cold Atoms
  - **Optical lattice**: Hubbard model with internal degrees of freedom

“Old” Ideas (and applications):

- Synthetic gauge field: [Jaksch and Zoller proposal and extensions](#)

New use:

- Models with **extra-dimension** (in particular 4D)

# Outline:

- Idea of Quantum Simulator ([Feynman](#))
- Cold Atoms
  - **Optical lattice**: Hubbard model with internal degrees of freedom

“Old” Ideas (and applications):

- Synthetic gauge field: [Jaksch and Zoller proposal and extensions](#)

New use:

- Models with **extra-dimension** (in particular 4D)
- Models with twisted boundary conditions (in progress)

# Outline:

- Idea of Quantum Simulator ([Feynman](#))
- Cold Atoms
  - **Optical lattice**: Hubbard model with internal degrees of freedom

“Old” Ideas (and applications):

- Synthetic gauge field: [Jaksch and Zoller proposal and extensions](#)

New use:

- Models with **extra-dimension** (in particular 4D)
- Models with twisted boundary conditions (in progress)
- Prospects



# Why a Quantum Simulator?

(after discretization) In a quantum system dim.  $\mathcal{H}$  grows **exponentially** with  $V$

# Why a Quantum Simulator?

(after discretization) In a quantum system dim.  $\mathcal{H}$  grows **exponentially** with  $V$

Feynman (1983):

- Classical computer simulation takes exponential time
- Hypothetical quantum computer does not ( $\simeq$  polynomial)

# Why a Quantum Simulator?

(after discretization) In a quantum system dim.  $\mathcal{H}$  grows **exponentially** with  $V$

**Feynman (1983):**

- Classical computer simulation takes exponential time
- Hypothetical quantum computer does not ( $\simeq$  polynomial)

⇒ New field: Quantum computation

# Why a Quantum Simulator?

(after discretization) In a quantum system dim.  $\mathcal{H}$  grows **exponentially** with  $V$

**Feynman (1983):**

- Classical computer simulation takes exponential time
- Hypothetical quantum computer does not ( $\simeq$  polynomial)

⇒ New field: Quantum computation

Universal quantum computer still far away

# Why a Quantum Simulator?

(after discretization) In a quantum system dim.  $\mathcal{H}$  grows **exponentially** with  $V$

**Feynman (1983):**

- Classical computer simulation takes exponential time
- Hypothetical quantum computer does not ( $\simeq$  polynomial)

⇒ New field: Quantum computation

Universal quantum computer still far away/**Dedicated quantum simulator possible**

# Why a Quantum Simulator?

(after discretization) In a quantum system dim.  $\mathcal{H}$  grows **exponentially** with  $V$

**Feynman (1983):**

- Classical computer simulation takes exponential time
- Hypothetical quantum computer does not ( $\simeq$  polynomial)

⇒ New field: Quantum computation

Universal quantum computer still far away/**Dedicated quantum simulator possible**

**Requirements:**

- reproduces  $\mathcal{H}_{eff}(\lambda_i)$
- parameters  $\lambda_i$  tunable

# Why a Quantum Simulator?

(after discretization) In a quantum system dim.  $\mathcal{H}$  grows **exponentially** with  $V$

**Feynman (1983):**

- Classical computer simulation takes exponential time
- Hypothetical quantum computer does not ( $\simeq$  polynomial)

⇒ New field: Quantum computation

Universal quantum computer still far away/**Dedicated quantum simulator possible**

**Requirements:**

- reproduces  $\mathcal{H}_{eff}(\lambda_i)$  **Analogic** vs **Digital**: reproduces  $e^{-i\mathcal{H}_{eff}(\lambda_i)t}$
- parameters  $\lambda_i$  tunable

# Why a Quantum Simulator?

(after discretization) In a quantum system dim.  $\mathcal{H}$  grows **exponentially** with  $V$

**Feynman (1983):**

- Classical computer simulation takes exponential time
- Hypothetical quantum computer does not ( $\simeq$  polynomial)

⇒ New field: Quantum computation

Universal quantum computer still far away/**Dedicated quantum simulator possible**

**Requirements:**

- reproduces  $\mathcal{H}_{eff}(\lambda_i)$  **Analogic** vs **Digital**: reproduces  $e^{-i\mathcal{H}_{eff}(\lambda_i)t}$
- parameters  $\lambda_i$  tunable

→ Good candidate: **Cold atoms**



# Untractable vs Unaccessible

Two lines of interest in Quantum simulation

# Untractable vs Unaccessible

Two lines of interest in Quantum simulation

- Solution of non-calculable models

# Untractable vs Unaccessible

Two lines of interest in Quantum simulation

- Solution of non-calculable models

Examples:

- Superconductivity
- (Non-abelian lattice) Gauge theory
- Non-equilibrium dynamics and time evolution
- ...

# Untractable vs Unaccessible

Two lines of interest in Quantum simulation

- Solution of non-calculable models
  - Examples:
    - Superconductivity
    - (Non-abelian lattice) Gauge theory
    - Non-equilibrium dynamics and time evolution
    - ...
- Design of systems unobservable (or yet unobserved)

# Untractable vs Unaccessible

Two lines of interest in Quantum simulation

- Solution of non-calculable models

Examples:

- Superconductivity
- (Non-abelian lattice) Gauge theory
- Non-equilibrium dynamics and time evolution
- ...

- Design of systems unobservable (or yet unobserved)

Examples:

- Zitterbewegung, Klein paradox
- Neutrino oscillations (direct measurement)
- Dirac fermion in artificially curved spacetime
- ...

# Untractable vs Unaccessible

## Two lines of interest in Quantum simulation

### ● Solution of non-calculable models

#### Examples:

- Superconductivity
- (Non-abelian lattice) Gauge theory
- Non-equilibrium dynamics and time evolution
- ...

### ● Design of systems unobservable (or yet unobserved)

#### Examples:

- Zitterbewegung, Klein paradox
- Neutrino oscillations (direct measurement)
- Dirac fermion in artificially curved spacetime
- ...
- Simulation of models with **four spatial dimensions**

# Untractable vs Unaccessible

Two lines of interest in Quantum simulation

- Solution of non-calculable models

Examples:

- Superconductivity
- (Non-abelian lattice) Gauge theory
- Non-equilibrium dynamics and time evolution
- ...

- Design of systems unobservable (or yet unobserved)

Examples:

- Zitterbewegung, Klein paradox
- Neutrino oscillations (direct measurement)
- Dirac fermion in artificially curved spacetime
- ...
- Simulation of models with **four spatial dimensions**

Note that once interactions are included the above problems become not calculable

# Observations & Ingredients for the simulation

- All the example mentioned can be conveniently formulated on the [lattice](#)



# Observations & Ingredients for the simulation

- All the example mentioned can be conveniently formulated on the **lattice**
  - Natural set-up: Optical lattices

# Observations & Ingredients for the simulation

- All the example mentioned can be conveniently formulated on the **lattice**
  - Natural set-up: Optical lattices
  
- The Hamiltonian to be simulated reduces to extended Hubbard model with internal degrees of freedom

# Observations & Ingredients for the simulation

- All the example mentioned can be conveniently formulated on the **lattice**
  - Natural set-up: Optical lattices
  
- The Hamiltonian to be simulated reduces to extended Hubbard model with internal degrees of freedom
  - Natural ingredients: Hyperfine states, Superlattice

# Observations & Ingredients for the simulation

- All the example mentioned can be conveniently formulated on the **lattice**
  - Natural set-up: Optical lattices
- The Hamiltonian to be simulated reduces to extended Hubbard model with internal degrees of freedom
  - Natural ingredients: Hyperfine states, Superlattice
- (Historical) Prototype: Simulation of Synthetic gauge field in optical lattices

# Observations & Ingredients for the simulation

- All the example mentioned can be conveniently formulated on the **lattice**
  - Natural set-up: Optical lattices
- The Hamiltonian to be simulated reduces to extended Hubbard model with internal degrees of freedom
  - Natural ingredients: Hyperfine states, Superlattice
- (Historical) Prototype: Simulation of Synthetic gauge field in optical lattices
  - We borrow some ideas from there

# Synthetic (external) gauge field

Motivation: Interesting **strong coupling** phenomena where the **charge** of electron is **important**

# Synthetic (external) gauge field

Motivation: Interesting **strong coupling** phenomena where the **charge** of electron is **important**

Examples:

- **Quantum Hall effect**: a constant **strong** magnetic field needed

# Synthetic (external) gauge field

Motivation: Interesting **strong coupling** phenomena where the **charge** of electron is **important**

Examples:

- **Quantum Hall effect**: a constant **strong** magnetic field needed
- More ambitious: **Chromodynamics**



# Synthetic (external) gauge field

Motivation: Interesting **strong coupling** phenomena where the **charge** of electron is **important**

Examples:

- **Quantum Hall effect**: a constant **strong** magnetic field needed
- More ambitious: **Chromodynamics**

**Problem**: atoms are **neutral**

# Synthetic (external) gauge field

Motivation: Interesting **strong coupling** phenomena where the **charge** of electron is **important**

Examples:

- **Quantum Hall effect**: a constant **strong** magnetic field needed
- More ambitious: **Chromodynamics**

**Problem**: atoms are **neutral**

- Old Idea: **rotation**  $\equiv$  **constant magnetic field** (Larmor theorem)

# Synthetic (external) gauge field

Motivation: Interesting **strong coupling** phenomena where the **charge** of electron is **important**

Examples:

- **Quantum Hall effect**: a constant **strong** magnetic field needed
- More ambitious: **Chromodynamics**

**Problem**: atoms are **neutral**

- Old Idea: **rotation**  $\equiv$  **constant magnetic field** (Larmor theorem)

 magnetic flux limited

# Synthetic (external) gauge field

Motivation: Interesting **strong coupling** phenomena where the **charge** of electron is **important**

Examples:

- **Quantum Hall effect**: a constant **strong** magnetic field needed
- More ambitious: **Chromodynamics**

**Problem**: atoms are **neutral**

- Old Idea: **rotation**  $\equiv$  **constant magnetic field** (Larmor theorem)
  - ☹ magnetic flux limited
- **Jaksch and Zoller, 2003** (theory): **artificial** magnetic field with **Raman lasers** in OL

# Synthetic (external) gauge field

Motivation: Interesting **strong coupling** phenomena where the **charge** of electron is **important**

Examples:

- **Quantum Hall effect**: a constant **strong** magnetic field needed
- More ambitious: **Chromodynamics**

**Problem**: atoms are **neutral**

- Old Idea: **rotation**  $\equiv$  **constant magnetic field** (Larmor theorem)
  - ☹ magnetic flux limited
- **Jaksch and Zoller, 2003** (theory): **artificial** magnetic field with **Raman lasers** in OL
  - access the **Hofstadter butterfly**'s physics

# Synthetic (external) gauge field

Motivation: Interesting **strong coupling** phenomena where the **charge** of electron is **important**

Examples:

- **Quantum Hall effect**: a constant **strong** magnetic field needed
- More ambitious: **Chromodynamics**

**Problem**: atoms are **neutral**

- Old Idea: **rotation**  $\equiv$  **constant magnetic field** (Larmor theorem)
  - ☹ magnetic flux limited
- **Jaksch and Zoller, 2003** (theory): **artificial** magnetic field with **Raman lasers** in OL
  - access the **Hofstadter butterfly**'s physics

Quickly developing research area

# Gauging the Hubbard model

**Example:** Electron (without spin) in a **2d** crystal, Hubbard model (small filling)

$$H = -J_x \sum_{m,n=-\infty}^{\infty} a_{m+1,n}^\dagger a_{m,n} - J_y \sum_{m,n=-\infty}^{\infty} a_{m,n+1}^\dagger a_{m,n} + h.c.$$

# Gauging the Hubbard model

**Example:** Electron (without spin) in a **2d** crystal, Hubbard model (small filling)

$$H = -J_x \sum_{m,n=-\infty}^{\infty} a_{m+1,n}^\dagger a_{m,n} - J_y \sum_{m,n=-\infty}^{\infty} a_{m,n+1}^\dagger a_{m,n} + h.c.$$

**Hofstadter,1976:** coupling to an external background field, **A vector potential**



# Gauging the Hubbard model

**Example:** Electron (without spin) in a **2d** crystal, Hubbard model (small filling)

$$H = -J_x \sum_{m,n=-\infty}^{\infty} a_{m+1,n}^\dagger a_{m,n} - J_y \sum_{m,n=-\infty}^{\infty} a_{m,n+1}^\dagger a_{m,n} + h.c.$$

**Hofstadter,1976:** coupling to an external background field, **A** vector potential

**Peierls substitution:** pseudo-momentum **k**  $\rightarrow$  **k** - **A**

# Gauging the Hubbard model

**Example:** Electron (without spin) in a **2d** crystal, Hubbard model (small filling)

$$H = -J_x \sum_{m,n=-\infty}^{\infty} a_{m+1,n}^\dagger a_{m,n} - J_y \sum_{m,n=-\infty}^{\infty} a_{m,n+1}^\dagger a_{m,n} + h.c.$$

**Hofstadter, 1976:** coupling to an external background field, **A** vector potential

**Peierls substitution:** pseudo-momentum  $\mathbf{k} \rightarrow \mathbf{k} - \mathbf{A}$

→ Phase changing hopping ( $\equiv$  **Wilson line**)

$$H = -J_x \sum_{m,n=-\infty}^{\infty} e^{iA_x(m,n)} a_{m+1,n}^\dagger a_{m,n} - J_y \sum_{m,n=-\infty}^{\infty} e^{iA_y(m,n)} a_{m,n+1}^\dagger a_{m,n} + h.c.$$

# Gauging the Hubbard model

**Example:** Electron (without spin) in a **2d** crystal, Hubbard model (small filling)

$$H = -J_x \sum_{m,n=-\infty}^{\infty} a_{m+1,n}^\dagger a_{m,n} - J_y \sum_{m,n=-\infty}^{\infty} a_{m,n+1}^\dagger a_{m,n} + h.c.$$

**Hofstadter, 1976:** coupling to an external background field, **A** vector potential

**Peierls substitution:** pseudo-momentum  $\mathbf{k} \rightarrow \mathbf{k} - \mathbf{A}$

$\rightarrow$  Phase changing hopping ( $\equiv$  **Wilson line**)

$$H = -J_x \sum_{m,n=-\infty}^{\infty} e^{iA_x(m,n)} a_{m+1,n}^\dagger a_{m,n} - J_y \sum_{m,n=-\infty}^{\infty} e^{iA_y(m,n)} a_{m,n+1}^\dagger a_{m,n} + h.c.$$

Constant magnetic field  $\mathbf{B} = (0, 0, B) \xrightarrow{\text{Landau gauge}} \mathbf{A} = (0, Bx, 0),$

# Gauging the Hubbard model

**Example:** Electron (without spin) in a **2d** crystal, Hubbard model (small filling)

$$H = -J_x \sum_{m,n=-\infty}^{\infty} a_{m+1,n}^\dagger a_{m,n} - J_y \sum_{m,n=-\infty}^{\infty} a_{m,n+1}^\dagger a_{m,n} + h.c.$$

**Hofstadter, 1976:** coupling to an external background field, **A** vector potential

**Peierls substitution:** pseudo-momentum  $\mathbf{k} \rightarrow \mathbf{k} - \mathbf{A}$

$\rightarrow$  Phase changing hopping ( $\equiv$  **Wilson line**)

$$H = -J_x \sum_{m,n=-\infty}^{\infty} e^{iA_x(m,n)} a_{m+1,n}^\dagger a_{m,n} - J_y \sum_{m,n=-\infty}^{\infty} e^{iA_y(m,n)} a_{m,n+1}^\dagger a_{m,n} + h.c.$$

Constant magnetic field  $\mathbf{B} = (0, 0, B) \xrightarrow{\text{Landau gauge}} \mathbf{A} = (0, Bx, 0),$

# Gauging the Hubbard model

**Example:** Electron (without spin) in a **2d** crystal, Hubbard model (small filling)

$$H = -J_x \sum_{m,n=-\infty}^{\infty} a_{m+1,n}^\dagger a_{m,n} - J_y \sum_{m,n=-\infty}^{\infty} a_{m,n+1}^\dagger a_{m,n} + h.c.$$

**Hofstadter, 1976:** coupling to an external background field, **A** vector potential

**Peierls substitution:** pseudo-momentum  $\mathbf{k} \rightarrow \mathbf{k} - \mathbf{A}$

$\rightarrow$  Phase changing hopping ( $\equiv$  **Wilson line**)

$$H = -J_x \sum_{m,n=-\infty}^{\infty} e^{iA_x(m,n)} a_{m+1,n}^\dagger a_{m,n} - J_y \sum_{m,n=-\infty}^{\infty} e^{iA_y(m,n)} a_{m,n+1}^\dagger a_{m,n} + h.c.$$

Constant magnetic field  $\mathbf{B} = (0, 0, B) \xrightarrow{\text{Landau gauge}} \mathbf{A} = (0, Bx, 0)$ ,

Generalization to **non abelian** theory in **D** dim.:

•  $a_{m,n} \rightarrow a_{\mathbf{r}}^\sigma$ ,  $\mathbf{r} \in \mathbb{R}^D$ , ex.  $SU(N)$ ,  $\sigma = 1, \dots, N$

•  $A_\mu \equiv A_\mu^I T_I$ ,  $T_I$  gauge g. generators, hopping **phase**  $\rightarrow$  hopping **matrix**

# Jaksch-Zoller (JZ) proposal

Two ingredients

# Jaksch-Zoller (JZ) proposal

## Two ingredients

- Superposition of two 2D lattices loaded with two different hyperfine states  $|g\rangle$  and  $|e\rangle$

# Jaksch-Zoller (JZ) proposal

## Two ingredients

- Superposition of two 2D lattices loaded with two different hyperfine states  $|g\rangle$  and  $|e\rangle$   
Off-set of  $\lambda/4$  in  $y$  direction,  $x_m = \lambda/2 m$ ,  $y_n = \lambda/4 n$



# Jaksch-Zoller (JZ) proposal

## Two ingredients

- Superposition of two 2D lattices loaded with two different hyperfine states  $|g\rangle$  and  $|e\rangle$   
Off-set of  $\lambda/4$  in  $y$  direction,  $x_m = \lambda/2 m$ ,  $y_n = \lambda/4 n$ ,  $V_{0y} \gg V_{0x} \Rightarrow J_{0y}$  suppressed

# Jaksch-Zoller (JZ) proposal

## Two ingredients

- Superposition of two 2D lattices loaded with two different hyperfine states  $|g\rangle$  and  $|e\rangle$   
Off-set of  $\lambda/4$  in  $y$  direction,  $x_m = \lambda/2 m$ ,  $y_n = \lambda/4 n$ ,  $V_{0y} \gg V_{0x} \Rightarrow J_{0y}$  suppressed
- Stimulated Rabi osc. between  $|g\rangle$  and  $|e\rangle$  by two running Raman lasers:  
$$\Omega(\mathbf{x}) = \Omega_0 e^{i(\mathbf{k}_e - \mathbf{k}_g) \cdot \mathbf{x}}$$

# Jaksch-Zoller (JZ) proposal

## Two ingredients

- Superposition of two 2D lattices loaded with two different hyperfine states  $|g\rangle$  and  $|e\rangle$   
Off-set of  $\lambda/4$  in  $y$  direction,  $x_m = \lambda/2 m$ ,  $y_n = \lambda/4 n$ ,  $V_{0y} \gg V_{0x} \Rightarrow J_{0y}$  suppressed
- Stimulated Rabi osc. between  $|g\rangle$  and  $|e\rangle$  by two running Raman lasers:  
 $\Omega(\mathbf{x}) = \Omega_0 e^{i(\mathbf{k}_e - \mathbf{k}_g) \cdot \mathbf{x}}$ , perturbation of OL (same Wannier),  $|\mathbf{k}_g| \simeq |\mathbf{k}_e| = \frac{2\pi}{\lambda}$

# Jaksch-Zoller (JZ) proposal

## Two ingredients

- Superposition of two 2D lattices loaded with two different hyperfine states  $|g\rangle$  and  $|e\rangle$   
Off-set of  $\lambda/4$  in  $y$  direction,  $x_m = \lambda/2 m$ ,  $y_n = \lambda/4 n$ ,  $V_{0y} \gg V_{0x} \Rightarrow J_{0y}$  suppressed
- Stimulated Rabi osc. between  $|g\rangle$  and  $|e\rangle$  by two running Raman lasers:  
 $\Omega(\mathbf{x}) = \Omega_0 e^{i(\mathbf{k}_e - \mathbf{k}_g) \cdot \mathbf{x}}$ , perturbation of OL (same Wannier),  $|\mathbf{k}_g| \simeq |\mathbf{k}_e| = \frac{2\pi}{\lambda}$

## Final result

$$H = J_{0x} \sum_{m,n=-\infty}^{\infty} a_{m+1,n}^\dagger a_{m,n} + \sum_{m,n=-\infty}^{\infty} J_y e^{iA_y(m,n)} a_{m,n+1}^\dagger a_{m,n} + h.c.,$$

# Jaksch-Zoller (JZ) proposal

## Two ingredients

- Superposition of two 2D lattices loaded with two different hyperfine states  $|g\rangle$  and  $|e\rangle$   
Off-set of  $\lambda/4$  in  $y$  direction,  $x_m = \lambda/2 m$ ,  $y_n = \lambda/4 n$ ,  $V_{0y} \gg V_{0x} \Rightarrow J_{0y}$  suppressed
- Stimulated Rabi osc. between  $|g\rangle$  and  $|e\rangle$  by two running Raman lasers:  
 $\Omega(\mathbf{x}) = \Omega_0 e^{i(\mathbf{k}_e - \mathbf{k}_g) \cdot \mathbf{x}}$ , perturbation of OL (same Wannier),  $|\mathbf{k}_g| \simeq |\mathbf{k}_e| = \frac{2\pi}{\lambda}$

## Final result

$$H = J_{0x} \sum_{m,n=-\infty}^{\infty} a_{m+1,n}^\dagger a_{m,n} + \sum_{m,n=-\infty}^{\infty} J_y e^{iA_y(m,n)} a_{m,n+1}^\dagger a_{m,n} + h.c.,$$

Average over Wannier  $J_y e^{iA_y(m,n)} = \frac{\hbar}{2} \int d^2\mathbf{x} \mathbf{w}^*(\mathbf{x} - \mathbf{x}_{m,n}) \Omega \mathbf{w}(\mathbf{x} - \mathbf{x}_{m,n+1})$ .

# Jaksch-Zoller (JZ) proposal

## Two ingredients

- Superposition of two 2D lattices loaded with two different hyperfine states  $|g\rangle$  and  $|e\rangle$   
Off-set of  $\lambda/4$  in  $y$  direction,  $x_m = \lambda/2 m$ ,  $y_n = \lambda/4 n$ ,  $V_{0y} \gg V_{0x} \Rightarrow J_{0y}$  suppressed
- Stimulated Rabi osc. between  $|g\rangle$  and  $|e\rangle$  by two running Raman lasers:  
 $\Omega(\mathbf{x}) = \Omega_0 e^{i(\mathbf{k}_e - \mathbf{k}_g) \cdot \mathbf{x}}$ , perturbation of OL (same Wannier),  $|\mathbf{k}_g| \simeq |\mathbf{k}_e| = \frac{2\pi}{\lambda}$

## Final result

$$H = J_{0x} \sum_{m,n=-\infty}^{\infty} a_{m+1,n}^\dagger a_{m,n} + \sum_{m,n=-\infty}^{\infty} J_y e^{iA_y(m,n)} a_{m,n+1}^\dagger a_{m,n} + h.c.,$$

Average over Wannier  $J_y e^{iA_y(m,n)} = \frac{\hbar}{2} \int d^2\mathbf{x} \mathbf{w}^*(\mathbf{x} - \mathbf{x}_{m,n}) \Omega \mathbf{w}(\mathbf{x} - \mathbf{x}_{m,n+1})$ .

For  $\mathbf{k}_e - \mathbf{k}_g // \hat{x} \rightarrow A_y(m,n) = 2\pi\Phi m$ , Constant magnetic field in Landau gauge

# Comments...

JZ proposal: experiment [PRL 107 255301 \(2012\)](#) (I. Bloch group)

# Comments...

**JZ proposal:** experiment [PRL 107 255301 \(2012\)](#) (I. Bloch group)

Many more proposals with or without lattice:



# Comments...

**JZ proposal:** experiment [PRL 107 255301 \(2012\)](#) (I. Bloch group)

Many more proposals with or without lattice:

- Rotation (experiments since 2000)

# Comments...

**JZ proposal:** experiment [PRL 107 255301 \(2012\)](#) (I. Bloch group)

Many more proposals with or without lattice:

- Rotation (experiments since 2000)
- Dark state (Berry phase), [first experiments NIST, 2009-2010](#)

# Comments...

**JZ proposal:** experiment [PRL 107 255301 \(2012\)](#) (I. Bloch group)

Many more proposals with or without lattice:

- Rotation (experiments since 2000)
- Dark state (Berry phase), [first experiments NIST, 2009-2010](#)
- Superlattices, etc      EX. [Toolbox...](#) by Mazza *et al.*, New J. Phys. 14 (2012) 015007

# Comments...

**JZ proposal:** experiment [PRL 107 255301 \(2012\)](#) (I. Bloch group)

Many more proposals with or without lattice:

- Rotation (experiments since 2000)
- Dark state (Berry phase), [first experiments NIST, 2009-2010](#)
- Superlattices, etc      EX. [Toolbox...](#) by Mazza *et al.*, New J. Phys. 14 (2012) 015007
- Shaking      cf talk K. Sacha & K. Sengstock

# Comments...

**JZ proposal:** experiment [PRL 107 255301 \(2012\)](#) (I. Bloch group)

Many more proposals with or without lattice:

- Rotation (experiments since 2000)
- Dark state (Berry phase), [first experiments NIST, 2009-2010](#)
- Superlattices, etc      EX. [Toolbox...](#) by Mazza *et al.*, New J. Phys. 14 (2012) 015007
- Shaking      cf talk K. Sacha & K. Sengstock

[Osterloh \*et al.\*, Ruseckas \*et al.\*, 2005](#) [Hyperfine states](#) as internal gauge d.o.f  
⇒ [non abelian](#) interactions

# Comments...

**JZ proposal:** experiment [PRL 107 255301 \(2012\)](#) (I. Bloch group)

Many more proposals with or without lattice:

- Rotation (experiments since 2000)
- Dark state (Berry phase), [first experiments NIST, 2009-2010](#)
- Superlattices, etc      EX. [Toolbox...](#) by Mazza *et al.*, New J. Phys. 14 (2012) 015007
- Shaking      cf talk K. Sacha & K. Sengstock

[Osterloh \*et al.\*, Ruseckas \*et al.\*, 2005](#) [Hyperfine states](#) as internal gauge d.o.f

⇒ [non abelian](#) interactions Ex: SU(2), hopping phase →  $2 \times 2$  unitary hopping matrices

# Comments...

**JZ proposal:** experiment [PRL 107 255301 \(2012\)](#) (I. Bloch group)

Many more proposals with or without lattice:

- Rotation (experiments since 2000)
- Dark state (Berry phase), [first experiments NIST, 2009-2010](#)
- Superlattices, etc      EX. [Toolbox...](#) by Mazza *et al.*, New J. Phys. 14 (2012) 015007
- Shaking    very recent Hamburg-Icfo-Dresden [non-abelian PRL 110, xxxxxx \(2012\)](#)

[Osterloh \*et al.\*, Ruseckas \*et al.\*, 2005](#) [Hyperfine states](#) as internal gauge d.o.f

⇒ [non abelian](#) interactions Ex: SU(2), hopping phase →  $2 \times 2$  unitary hopping matrices

# Comments...

**JZ proposal:** experiment [PRL 107 255301 \(2012\)](#) (I. Bloch group)

Many more proposals with or without lattice:

- Rotation (experiments since 2000)
- Dark state (Berry phase), [first experiments NIST, 2009-2010](#)
- Superlattices, etc      EX. [Toolbox...](#) by Mazza *et al.*, New J. Phys. 14 (2012) 015007
- Shaking    very recent Hamburg-Icfo-Dresden [non-abelian PRL 110, xxxxxx \(2012\)](#)

[Osterloh \*et al.\*, Ruseckas \*et al.\*, 2005](#) [Hyperfine states](#) as internal gauge d.o.f

⇒ [non abelian](#) interactions Ex: SU(2), hopping phase →  $2 \times 2$  unitary hopping matrices

Observations:

- classical background gauge field configuration (no dynamics)
- [always gauge-fixed Hamiltonian](#)
- relativistic matter & [extradimension](#) simulations require synthetic gauge field



# Modelling extra-dimensions (Boada,AC,Latorre,Lewenstein, PRL 2012)

(during the workshop) in OL 3D Hubbard model  $\rightarrow$  1D,2D by tuning optical potential

# Modelling extra-dimensions (Boada,AC,Latorre,Lewenstein, PRL 2012)

(during the workshop) in OL 3D Hubbard model  $\rightarrow$  1D,2D by tuning optical potential and  $>$  3D?

# Modelling extra-dimensions (Boada,AC,Latorre,Lewenstein, PRL 2012)

(during the workshop) in OL 3D Hubbard model  $\rightarrow$  1D,2D by tuning optical potential and  $>$  3D? In a lattice **Dimensionality**  $\equiv$  **Connectivity**

# Modelling extra-dimensions (Boada,AC,Latorre,Lewenstein, PRL 2012)

(during the workshop) in OL 3D Hubbard model  $\rightarrow$  1D,2D by tuning optical potential and  $> 3D$ ? In a lattice **Dimensionality**  $\equiv$  **Connectivity**

Basic idea: just kinetic term in  $D+1$  hypercubic lattice

$$H = -J \sum_{\mathbf{q}} \sum_{j=1}^{D+1} a_{\mathbf{q}+\mathbf{u}_j}^\dagger a_{\mathbf{q}} + H.c. ,$$

# Modelling extra-dimensions (Boada,AC,Latorre,Lewenstein, PRL 2012)

(during the workshop) in OL 3D Hubbard model  $\rightarrow$  1D,2D by tuning optical potential and  $> 3D$ ? In a lattice **Dimensionality**  $\equiv$  **Connectivity**

Basic idea: just kinetic term in  $D+1$  hypercubic lattice

$$H = -J \sum_{\mathbf{q}} \sum_{j=1}^{D+1} a_{\mathbf{q}+\mathbf{u}_j}^\dagger a_{\mathbf{q}} + H.c. ,$$

$\mathbf{q} = (\mathbf{r}, \sigma)$ ,  $D + 1$ -space in  $D$ -layers

$$H = -J \sum_{\mathbf{r}, \sigma} \left( \sum_{j=1}^D a_{\mathbf{r}+\mathbf{u}_j}^{(\sigma)\dagger} a_{\mathbf{r}}^{(\sigma)} + a_{\mathbf{r}}^{(\sigma+1)\dagger} a_{\mathbf{r}}^{(\sigma)} \right) + H.c. ,$$

# Modelling extra-dimensions (Boada,AC,Latorre,Lewenstein, PRL 2012)

(during the workshop) in OL 3D Hubbard model  $\rightarrow$  1D,2D by tuning optical potential and  $>$  3D? In a lattice **Dimensionality**  $\equiv$  **Connectivity**

Basic idea: just kinetic term in  $D+1$  hypercubic lattice

$$H = -J \sum_{\mathbf{q}} \sum_{j=1}^{D+1} a_{\mathbf{q}+\mathbf{u}_j}^\dagger a_{\mathbf{q}} + H.c. ,$$

$\mathbf{q} = (\mathbf{r}, \sigma)$ ,  $D + 1$ -space in  $D$ -layers

$$H = -J \sum_{\mathbf{r}, \sigma} \left( \sum_{j=1}^D a_{\mathbf{r}+\mathbf{u}_j}^{(\sigma)\dagger} a_{\mathbf{r}}^{(\sigma)} + a_{\mathbf{r}}^{(\sigma+1)\dagger} a_{\mathbf{r}}^{(\sigma)} \right) + H.c. ,$$

$a_{\mathbf{r}}^{(\sigma)}, a_{\mathbf{r}}^{(\sigma)\dagger}$  as Fock generators of different species on same site, **extra-dim**  $\rightarrow$  **internal d.o.f**

Hopping in extra-dim  $\rightarrow$  transmutation between “contiguous” species,

# Modelling extra-dimensions (Boada,AC,Latorre,Lewenstein, PRL 2012)

(during the workshop) in OL 3D Hubbard model  $\rightarrow$  1D,2D by tuning optical potential and  $>$  3D? In a lattice **Dimensionality**  $\equiv$  **Connectivity**

Basic idea: just kinetic term in  $D+1$  hypercubic lattice

$$H = -J \sum_{\mathbf{q}} \sum_{j=1}^{D+1} a_{\mathbf{q}+\mathbf{u}_j}^\dagger a_{\mathbf{q}} + H.c. ,$$

$\mathbf{q} = (\mathbf{r}, \sigma)$ ,  $D + 1$ -space in  $D$ -layers

$$H = -J \sum_{\mathbf{r}, \sigma} \left( \sum_{j=1}^D a_{\mathbf{r}+\mathbf{u}_j}^{(\sigma)\dagger} a_{\mathbf{r}}^{(\sigma)} + a_{\mathbf{r}}^{(\sigma+1)\dagger} a_{\mathbf{r}}^{(\sigma)} \right) + H.c. ,$$

$a_{\mathbf{r}}^{(\sigma)}, a_{\mathbf{r}}^{(\sigma)\dagger}$  as Fock generators of different species on same site, **extra-dim**  $\rightarrow$  **internal d.o.f**  
Hopping in extra-dim  $\rightarrow$  transmutation between “contiguous” species, 2 ways:

- State-dependent lattice (cf JZ proposal)
- On-site dressed lattice (cf relativistic fermions Toolbox)

# Extra-dimension implementation & detection

State-dependent & On-site dressed, common features

- Open or **periodic** boundary cond. in the extra-dim → **Compactification**
- **Complex** hopping on the  $S^1$  (p.b.c.) → **Flux** compactification
- Interactions (on-site and long range)



# Extra-dimension implementation & detection

State-dependent & On-site dressed, common features

- Open or **periodic** boundary cond. in the extra-dim → **Compactification**
- **Complex** hopping on the  $S^1$  (p.b.c.) → **Flux** compactification
- Interactions (on-site and long range)

State-dependent & On-site dressed, specific features

# Extra-dimension implementation & detection

State-dependent & On-site dressed, common features

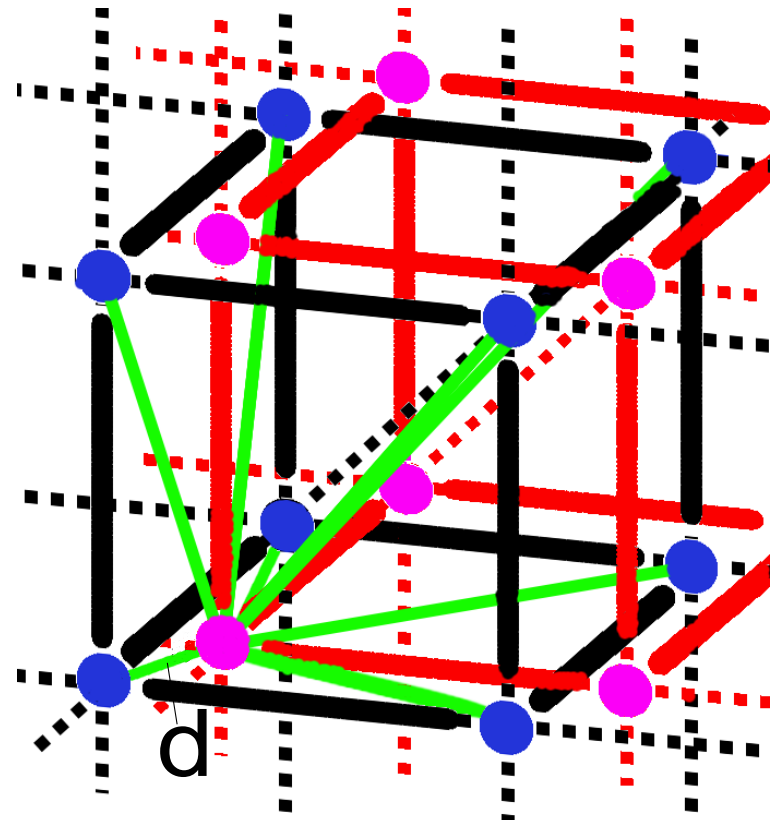
- Open or **periodic** boundary cond. in the extra-dim → **Compactification**
- **Complex** hopping on the  $S^1$  (p.b.c.) → **Flux** compactification
- Interactions (on-site and long range)

State-dependent & On-site dressed, specific features

- State-dependent lattice: easy realization

**Example:**  $^{87}\text{Rb}$ ,  $F = 1$ ,  $F = 2 \rightarrow$

**Bivolume**  $\equiv N = 2$  layers



# Extra-dimension implementation & detection

State-dependent & On-site dressed, common features

- Open or **periodic** boundary cond. in the extra-dim → **Compactification**
- **Complex** hopping on the  $S^1$  (p.b.c.) → **Flux** compactification
- Interactions (on-site and long range)

State-dependent & On-site dressed, specific features

- State-dependent lattice: easy realization  
**Example:**  $^{87}\text{Rb}$ ,  $F = 1$ ,  $F = 2$  → **Bivolume**  $\equiv N = 2$  layers
- On-site dressed lattice: up to  $N = 10$ , (alkali-earth  $\text{SU}(N)$  inv. interaction)  
**Example:**  $^{87}\text{Sr}$ ,  $^{40}\text{K}$ ,  $F = \frac{9}{2}$  (fermions)

# Extra-dimension implementation & detection

State-dependent & On-site dressed, common features

- Open or **periodic** boundary cond. in the extra-dim → **Compactification**
- **Complex** hopping on the  $S^1$  (p.b.c.) → **Flux** compactification
- Interactions (on-site and long range)

State-dependent & On-site dressed, specific features

- State-dependent lattice: easy realization  
**Example:**  $^{87}\text{Rb}$ ,  $F = 1$ ,  $F = 2$  → **Bivolume**  $\equiv N = 2$  layers
- On-site dressed lattice: up to  $N = 10$ , (alkali-earth  $\text{SU}(N)$  inv. interaction)  
**Example:**  $^{87}\text{Sr}$ ,  $^{40}\text{K}$ ,  $F = \frac{9}{2}$  (fermions), ?  $^{167}\text{Er}$ ,  $F = \frac{19}{2}$ ,  $N = 20$  ?

# Extra-dimension implementation & detection

State-dependent & On-site dressed, common features

- Open or **periodic** boundary cond. in the extra-dim → **Compactification**
- **Complex** hopping on the  $S^1$  (p.b.c.) → **Flux** compactification
- Interactions (on-site and long range)

State-dependent & On-site dressed, specific features

- State-dependent lattice: easy realization  
**Example:**  $^{87}\text{Rb}$ ,  $F = 1$ ,  $F = 2$  → **Bivolume**  $\equiv N = 2$  layers
- On-site dressed lattice: up to  $N = 10$ , (alkali-earth  $\text{SU}(N)$  inv. interaction)  
**Example:**  $^{87}\text{Sr}$ ,  $^{40}\text{K}$ ,  $F = \frac{9}{2}$  (fermions), ?  $^{167}\text{Er}$ ,  $F = \frac{19}{2}$ ,  $N = 20$  ?

Observables: dimensionality dependence

- Single-particle, scaling properties, **Example:** density of states
- Many-body, phase diagram, **Example:** MI-Superfluid

→ conclusion

# Exotic Boundary Conditions same people + J.Rodriguez-Laguna

Usual optical lattices: only **Open** Boundary Condition (BC) are possible

# Exotic Boundary Conditions same people + J.Rodriguez-Laguna

Usual optical lattices: only **Open** Boundary Condition (BC) are possible

In lattices with “compact” **extradimension** or Bessel lattices: **Periodic** BC

# Exotic Boundary Conditions same people + J.Rodriguez-Laguna

Usual optical lattices: only **Open** Boundary Condition (BC) are possible

In lattices with “compact” **extradimension** or Bessel lattices: **Periodic** BC

Can we do more?



# Exotic Boundary Conditions same people + J.Rodriguez-Laguna

Usual optical lattices: only **Open** Boundary Condition (BC) are possible

In lattices with “compact” **extradimension** or Bessel lattices: **Periodic** BC

Can we do more? For instance, **Twisted** BC  $\rightarrow$  Möbius strip or Klein bottle

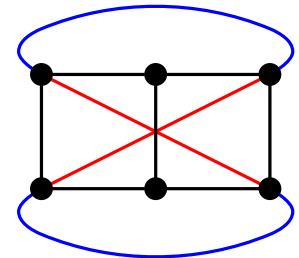
# Exotic Boundary Conditions same people + J.Rodriguez-Laguna

Usual optical lattices: only **Open** Boundary Condition (BC) are possible

In lattices with “compact” **extradimension** or Bessel lattices: **Periodic** BC

Can we do more? For instance, **Twisted** BC  $\rightarrow$  Möbius strip or Klein bottle

- Simplest way: construct the full lattice with internal states



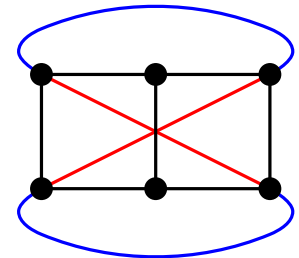
# Exotic Boundary Conditions same people + J.Rodriguez-Laguna

Usual optical lattices: only **Open** Boundary Condition (BC) are possible

In lattices with “compact” **extradimension** or Bessel lattices: **Periodic** BC

Can we do more? For instance, **Twisted** BC  $\rightarrow$  Möbius strip or Klein bottle

- Simplest way: construct the full lattice with internal states
  - Limitation: dimension of the lattice
  - Advantage: many copies of it



# Exotic Boundary Conditions same people + J.Rodriguez-Laguna

Usual optical lattices: only **Open** Boundary Condition (BC) are possible

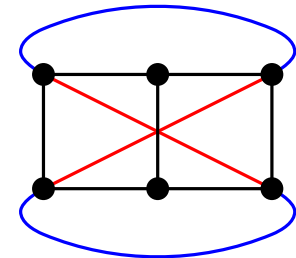
In lattices with “compact” **extradimension** or Bessel lattices: **Periodic** BC

Can we do more? For instance, **Twisted** BC  $\rightarrow$  Möbius strip or Klein bottle

- Simplest way: construct the full lattice with internal states

- Limitation: dimension of the lattice

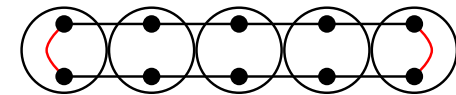
- Advantage: many copies of it



- Alternative: mixing of internal and spatial d.o.f

Idea: Two internal states  $+$  **Single site addressing**

$\rightarrow$  **Scalable circle**



# Exotic Boundary Conditions same people + J.Rodriguez-Laguna

Usual optical lattices: only **Open** Boundary Condition (BC) are possible

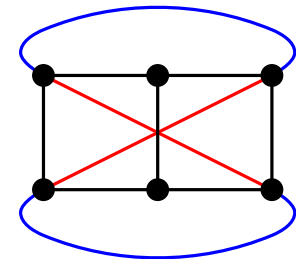
In lattices with “compact” **extradimension** or Bessel lattices: **Periodic** BC

Can we do more? For instance, **Twisted** BC  $\rightarrow$  Möbius strip or Klein bottle

● Simplest way: construct the full lattice with internal states

● Limitation: dimension of the lattice

● Advantage: many copies of it



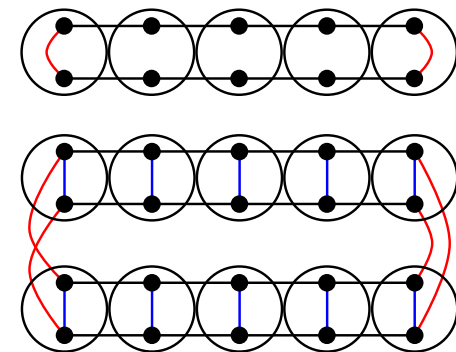
● Alternative: mixing of internal and spatial d.o.f

Idea: Two internal states  $+$  **Single site addressing**

$\rightarrow$  **Scalable circle**

Four internal states  $+$  **Single site addressing**

$\rightarrow$  **Scalable Möbius strip**



# Exotic Boundary Conditions same people + J.Rodriguez-Laguna

Usual optical lattices: only **Open** Boundary Condition (BC) are possible

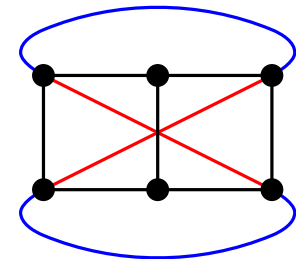
In lattices with “compact” **extradimension** or Bessel lattices: **Periodic** BC

Can we do more? For instance, **Twisted** BC  $\rightarrow$  Möbius strip or Klein bottle

● Simplest way: construct the full lattice with internal states

● Limitation: dimension of the lattice

● Advantage: many copies of it



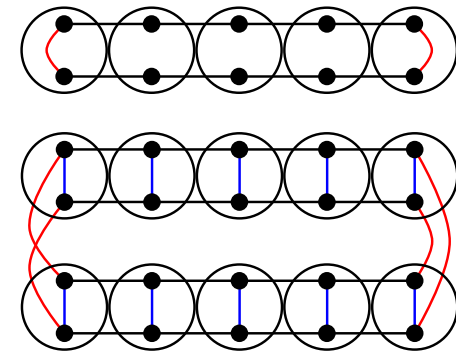
● Alternative: mixing of internal and spatial d.o.f

Idea: Two internal states  $+$  **Single site addressing**

$\rightarrow$  **Scalable circle**

Four internal states  $+$  **Single site addressing**

$\rightarrow$  **Scalable Möbius strip**



Interest: helicity modulus and stiffness in Heisenberg **[Fisher '71], [Rossini,'11]**

# Exotic Boundary Conditions same people + J.Rodriguez-Laguna

Usual optical lattices: only **Open** Boundary Condition (BC) are possible

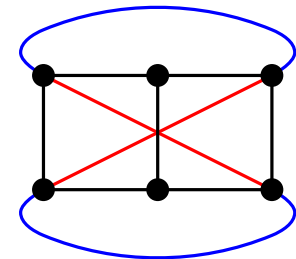
In lattices with “compact” **extradimension** or Bessel lattices: **Periodic** BC

Can we do more? For instance, **Twisted** BC  $\rightarrow$  Möbius strip or Klein bottle

- Simplest way: construct the full lattice with internal states

- Limitation: dimension of the lattice

- Advantage: many copies of it



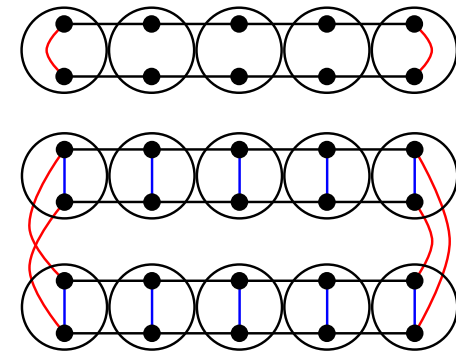
- Alternative: mixing of internal and spatial d.o.f

Idea: Two internal states  $+$  **Single site addressing**

$\rightarrow$  **Scalable circle**

Four internal states  $+$  **Single site addressing**

$\rightarrow$  **Scalable Möbius strip**



Interest: helicity modulus and stiffness in Heisenberg [**Fisher '71**], [**Rossini,'11**]

Additional ingredients: Generalized Twisted B.C. (with phases), Interactions

# Prospects



# Prospects

Very exciting time for Quantum simulation:

- Increasing number of theoretical proposals for [QFT simulators](#) in OL
- New experiments are coming

# Prospects

Very exciting time for Quantum simulation:

- Increasing number of theoretical proposals for **QFT simulators** in OL
- New experiments are coming

**My plans**, both in **Uncalculable**

- non-abelian link model (Gauge theory)

# Prospects

Very exciting time for Quantum simulation:

- Increasing number of theoretical proposals for **QFT simulators** in OL
- New experiments are coming

**My plans**, both in **Uncalculable**

- non-abelian link model (Gauge theory)

and **Unaccessible** problems:

# Prospects

Very exciting time for Quantum simulation:

- Increasing number of theoretical proposals for **QFT simulators** in OL
- New experiments are coming

**My plans**, both in **Uncalculable**

- non-abelian link model (Gauge theory)

and **Unaccessible** problems:

- Testing of Fulling-Unruh Thermalization theorem by simulation of artificial Dirac fermions in Rindler spacetime

# Prospects

Very exciting time for Quantum simulation:

- Increasing number of theoretical proposals for **QFT simulators** in OL
- New experiments are coming

**My plans**, both in **Uncalculable**

- non-abelian link model (Gauge theory)

and **Unaccessible** problems:

- Testing of Fulling-Unruh Thermalization theorem by simulation of artificial Dirac fermions in Rindler spacetime
- Simulation of Gravitino  $\equiv$  relativistic spin  $\frac{3}{2}$  particle

# Prospects

Very exciting time for Quantum simulation:

- Increasing number of theoretical proposals for **QFT simulators** in OL
- New experiments are coming

**My plans**, both in **Uncalculable**

- non-abelian link model (Gauge theory)

and **Unaccessible** problems:

- Testing of Fulling-Unruh Thermalization theorem by simulation of artificial Dirac fermions in Rindler spacetime
- Simulation of Gravitino  $\equiv$  relativistic spin  $\frac{3}{2}$  particle
- Extra-dimensions: exotic phases in Bosons (in progress) and Fermions

# Prospects

Very exciting time for Quantum simulation:

- Increasing number of theoretical proposals for **QFT simulators** in OL
- New experiments are coming

**My plans**, both in **Uncalculable**

- non-abelian link model (Gauge theory)

and **Unaccessible** problems:

- Testing of Fulling-Unruh Thermalization theorem by simulation of artificial Dirac fermions in Rindler spacetime
- Simulation of Gravitino  $\equiv$  relativistic spin  $\frac{3}{2}$  particle
- Extra-dimensions: exotic phases in Bosons (in progress) and Fermions
- Twisted Boundary Condition: robust effects in thermodynamics limits